

# LATTICE THERMAL EXPANSION OF LEAD IODIDE

LEAD iodide has a structure similar to that of cadmium iodide. Cadmium iodide is hexagonal with a typical 'layer structure', the interatomic forces within the layers being stronger than those perpendicular to the layers; the thermal expansion in a direction parallel to the layers is smaller than that in a perpendicular direction<sup>1</sup>. We have studied the thermal expansion behaviour of lead iodide by the X-ray method.

Analar grade  $\text{PbI}_2$  powder was used to obtain X-ray photographs at various temperatures from room temperature to  $114^\circ\text{C}$ . With filtered copper radiation, seven  $\text{CuK}\alpha_1\alpha_2$  doublets could be recorded on a flat-plate back-reflection camera. The experimental arrangement and the procedure for the evaluation of lattice parameters were the same as described earlier<sup>2</sup>. The values of the parameters at different temperatures are given in Table I. The lattice parameters at  $25^\circ\text{C}$  are  $a = 4.5562 \pm 0.0004 \text{ \AA}$  and  $c = 6.9830 \pm 0.0004 \text{ \AA}$ . These values agree well with those given by Swanson *et al.*<sup>3</sup>

TABLE I  
Lattice parameters of lead iodide

$t^\circ\text{C}$	$a \text{ \AA}$	$c \text{ \AA}$
25.0	4.5562	6.9830
64.5	4.5628	6.9905
86.0	4.5667	6.9964
104.0	4.5693	7.0017
114.0	4.5717	7.0035

The data on lattice parameters have been used to evaluate the directional coefficients of thermal expansion  $\bar{\alpha}_a = (1/a_{25})(\Delta a/\Delta t)$  and  $\bar{\alpha}_c = (1/c_{25})(\Delta c/\Delta t)$  over the range  $25\text{--}100^\circ\text{C}$ . The values are 36.6 and 32.3 (both in units of  $10^{-6}/^\circ\text{C}$ ). These values lead to  $105.5 \times 10^{-6}/^\circ\text{C}$  as the coefficient of cubical expansion. This value compares well with the values 100.8 and 108.6 (in units of  $10^{-6}/^\circ\text{C}$ ) reported by Fizeau<sup>4</sup> and Klemm *et al.*<sup>5</sup>.

We have also evaluated the coefficients defined by  $\alpha_a = (1/a_{25})(da/dt)$  and  $\alpha_c = (1/c_{25})(dc/dt)$  at various temperatures. These values are shown in Fig. 1.

The temperature variation of these coefficients can be represented by the equations:

$$\alpha_a = 26.36 \times 10^{-6} + 16.5 \times 10^{-8} t$$

$$\alpha_c = 15.69 \times 10^{-6} + 26.3 \times 10^{-8} t$$

where  $t$  is the temperature in  $^\circ\text{C}$ .

Regarding the anisotropy in thermal expansion, it is seen that the thermal expansion coefficient perpendicular to the layer ( $\alpha_c$ ), is smaller than that parallel to the layer ( $\alpha_a$ ). The anisotropy decreases as the temperature increases and at about  $100^\circ\text{C}$ , the two coefficients are nearly equal. If the trend continues, the anisotropy will be reversed at higher temperatures. With our technique, we could not record photographs of a measurable quality at higher temperatures.

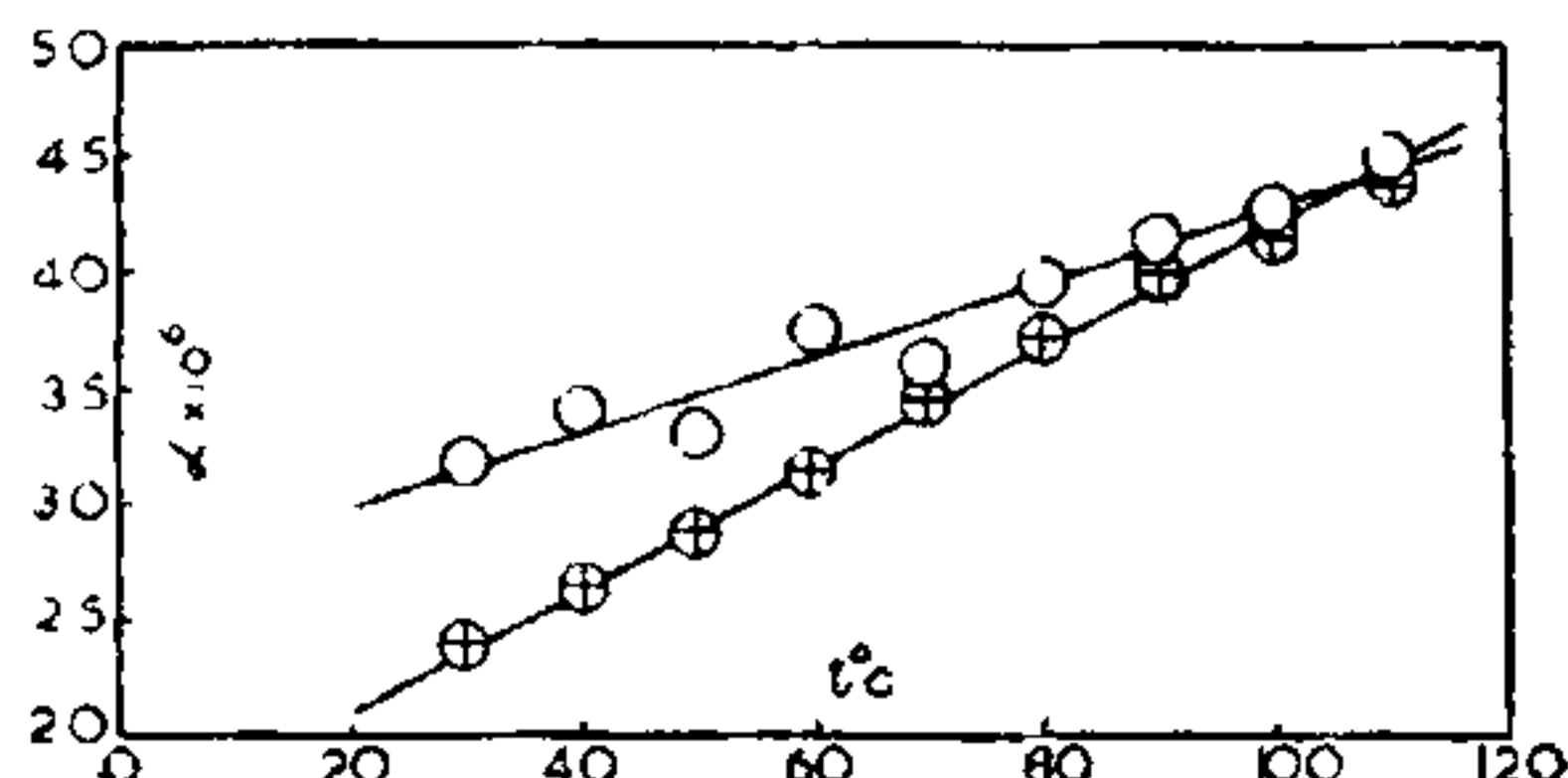


FIG. 1. Principal coefficients of expansion  $\alpha_a$  [O] and  $\alpha_c$  [⊕] of  $\text{PbI}_2$ .

In view of the similarity in the structure of  $\text{CdI}_2$  and  $\text{PbI}_2$ , we expected that the thermal expansion will have similar anisotropy. But at room temperature, the anisotropy in the thermal expansion of  $\text{PbI}_2$  is the opposite of that in the case of  $\text{CdI}_2$ . This, however, is not unusual. Wooster<sup>6</sup> and Bredig<sup>1</sup> have given examples of crystals with a layer structure having larger expansion coefficients parallel to the layer than perpendicular to the layer. A better understanding of the difference in the anisotropy of thermal expansion of  $\text{PbI}_2$  and  $\text{CdI}_2$  would be possible only when the elastic constants of these crystals become available.

Physics Department, D. B. SIRDESHMUKH,  
Post-Graduate Centre,  
Warangal (A.P.),

and  
Physics Department, V. T. DESHPANDE,  
Osmania University,  
Hyderabad-7, December 7, 1971.

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